


Search for Boosted Dark Matter Interacting with Electrons in Super-Kamiokande

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A search for boosted dark matter using 161.9 kt yr of Super-Kamiokande IV data is presented. We search for an excess of elastically scattered electrons above the atmospheric neutrino background, with a visible energy between 100 MeV and 1 TeV, pointing back to the Galactic center or the Sun. No such excess is observed. Limits on boosted dark matter event rates in multiple angular cones around the Galactic center and Sun are calculated. Limits are also calculated for a baseline model of boosted dark matter produced from cold dark matter annihilation or decay. This is the first experimental search for boosted dark matter from the Galactic center or the Sun interacting in a terrestrial detector.

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While there has long been ample evidence for the existence of dark matter [1–5], the specific properties and identity of dark matter remain elusive. The Λ CDM cosmology, which consists of long-lived dark matter that was nonrelativistic (“cold”) at freeze-out and a cosmological constant Λ , which corresponds to dark energy, has been well supported by cosmological observations [6]. Under this cosmology, the dark matter abundance has been measured by the observation of the cosmic microwave background to account for about 25% of the energy density of the Universe [7,8]. However, despite numerous direct and indirect detection searches, as well as searches for dark matter produced at particle accelerators, there has thus far been no definitive observation of particle dark matter [9–17].

With the properties of dark matter so uncertain, various possibilities must be considered. One possibility is that some dark matter is, in fact, not cold but is highly relativistic and has been produced at late times, thus denoted “boosted” dark matter [18–25]. Boosted dark matter could exist as a subdominant dark matter component, with a dominant cold

dark matter component accounting for most of the dark matter energy density of the Universe. In this way, boosted dark matter can remain consistent with Λ CDM. The subdominant boosted dark matter can be the same particle as the dominant cold dark matter, or it can be a different, lighter particle. Boosted dark matter can be produced from the dominant cold dark matter through a variety of processes, including annihilation [26,27], semiannihilation [28–32], number-changing $3 \rightarrow 2$ self-annihilation [33–35], and decay [22,36]. Boosted dark matter can then be observed through its scattering off electrons or nuclei in large-volume terrestrial detectors [37,38]. Current direct detection limits can be evaded in multicomponent models by having only the boosted dark matter species couple directly to standard model particles [18,22–24] or in boosted dark matter single-component models by invoking a spin-dependent dark matter-nucleon cross section [24].

This Letter reports the results of a search for boosted dark matter coupling to electrons in Super-Kamiokande (SK), with the boosted dark matter originating in the Galactic center or the Sun and with scattered electron energies ranging from 100 MeV to 1 TeV. This is the first time that this class of high-energy “electron elastic scatter-like” events has been studied at SK. The search is performed on 2628.1 days of SK-IV data, which corresponds to 161.9 kt yr exposure. The analysis is designed to be independent of the particular model of the coupling

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between boosted dark matter and electrons. This way, the results can be applied to any model that predicts a source of particles from the Galactic center or Sun which would scatter electrons to energies greater than 100 MeV.

The SK detector [39] is a water Cherenkov detector located 1000 m below Mt. Ikenoyama in Gifu, Japan. It consists of a 50 kt cylindrical tank of water, which is divided into a 32 kt (22.5 kt fiducial) inner detector (ID) surrounded by an outer detector (OD). The ID and the OD are optically separated by black Tyvek sheeting, and the ID is observed by 11 129 inward-facing 20-inch photomultiplier tubes (PMTs), while the OD is observed by 1885 outward-facing 8-inch PMTs. The ID provides most of the information used in event reconstruction, while the OD is used as an active veto region. Events are classified as fully contained (FC) if there is activity only in the ID. Relativistic charged particles in SK produce Cherenkov rings, which are categorized as e -like for electrons and γ 's or μ -like for muons.

The search begins with the fully contained fiducial-volume (FCFV) data set. This data set consists of all events with no OD activity, a reconstructed vertex inside the fiducial volume, and greater than 100 MeV of e -like momentum. These are events that originate and deposit all of their energy in the ID. This is a standard SK data set used to study atmospheric neutrinos [40]. From this data set, we search for elastically scattered electrons by applying the following analysis cuts: (i) one-ring (if $E_{\text{vis}} < 100$ GeV), (ii) e -like, (iii) zero decay electrons, and (iv) zero tagged neutrons. The first two cuts search for a single relativistic electron, while the final two cuts remove events with a signature of a nuclear interaction. Decay electrons in e -like events are the result of the $\pi^\pm \rightarrow \mu^\pm \rightarrow e^\pm$ decay chain with the π^\pm coming from a neutrino-nucleus interaction. Tagged neutrons originate from neutrons being knocked out of the nucleus following a neutrino-nucleus interaction, thermalizing, and capturing on hydrogen. Neutron captures are particularly numerous following neutrino deep inelastic scattering. Neither decay electrons nor neutron captures should occur following the elastic scatter of an electron by a boosted dark matter particle. The one-ring cut is not applied for events with visible energy above 100 GeV, as the ring-counting algorithm, which is tuned for lower-energy events, becomes unreliable at such high energies. We choose to restrict this analysis to SK-IV data only in order to take advantage of neutron tagging to remove atmospheric neutrino background. While neutron tagging has been used similarly in SK proton decay searches [41], this is the first time that it has been applied toward sample purification at energies greater than about 1 GeV. Event displays of example events passing the analysis cuts are provided in Supplemental Material [42].

Because of intricacies of the SK-IV trigger logic, events $\gtrsim 50$ GeV often cannot have neutron tagging applied.

Events for which neutron tagging cannot be applied are considered a signal if they pass the first three analysis cuts. Fourteen such events were found by FCFV selection, none of which passed the first three analysis cuts. This effect is accounted for in the signal efficiency, background rate, and limit calculations described below. The effect is minimal, since the efficiency of the neutron tagging cut is nearly 100%, and the background rate is driven down at such high energies by the sharp dropoff of the atmospheric neutrino flux.

The efficiency of the analysis cuts was found using a signal MC of 200 000 electrons with energies ranging from 30 MeV to 1 TeV. Events were simulated up to 1 m outside the fiducial volume, and efficiency was defined as the number of events passing each cut divided by the number of events simulated in the fiducial volume (events can migrate into or out of the fiducial volume during reconstruction). The efficiency of the FCFV selection and analysis cuts is shown in Fig. 1. The main cause of the reduction of efficiency with energy is the loss of containment at high energies; some higher-energy electromagnetic showers are able to penetrate from the FV into the OD and so do not pass the FC selection. The systematic uncertainty on the efficiency is estimated to be about 2%, most of which is due to uncertainty in the efficiency of the one-ring cut. The small size of the uncertainty of the FCFV selection and first three cuts is due to a significant separation between passing and failing events, so that very few signal events occur near the boundaries of the cuts. The efficiency of the neutron tagging cut is found by measuring the rate of false neutron

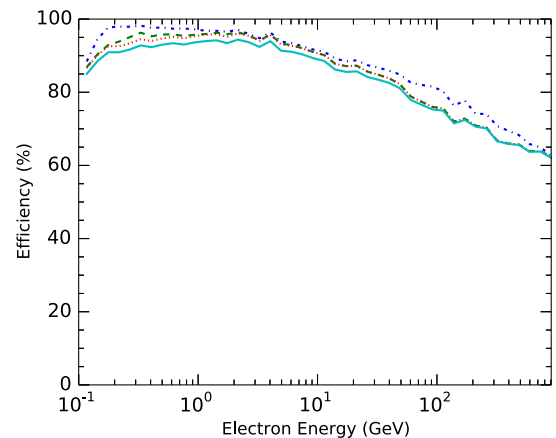


FIG. 1. Signal efficiency of the FCFV selection and analysis cuts as a function of the energy. Beginning with the FCFV selection (dashed-dotted blue curve), the addition of the one-ring (for $E_{\text{vis}} < 100$ GeV, dashed green curve), e -like (dotted red curve), and finally zero decay electrons and zero tagged neutrons cuts to arrive at the final efficiency (solid cyan curve) are shown. The efficiency of the zero decay electrons cut is $> 99.99\%$, so that the drop from the dotted red line to the solid cyan line is due solely to the neutron tagging cut.

TABLE I. Number of events over the entire sky passing each cut in 2628.1 days of SK4 data, simulated ν -MC background expectation, and signal efficiency at representative energy after each cut.

	100 MeV < E_{vis} < 1.33 GeV			1.33 GeV < E_{vis} < 20 GeV			$E_{\text{vis}} > 20$ GeV		
	Data	ν - MC	$\epsilon_{\text{sig}}(0.5 \text{ GeV})$	Data	ν - MC	$\epsilon_{\text{sig}}(5 \text{ GeV})$	Data	ν - MC	$\epsilon_{\text{sig}}(50 \text{ GeV})$
FCFV	15 206	14 858.1	97.7%	4908	5109.7	93.8%	118	107.5	84.9%
& single ring	11 367	10 997.4	95.8%	2868	3161.8	93.3%	71	68.2	82.2%
& e -like	5655	5571.5	94.7%	1514	1644.2	93.0%	71	68.1	82.2%
& zero decay- e	5049	5013.8	94.7%	1065	1207.2	93.0%	13	15.7	82.2%
& zero neutrons	4042	3992.9	93.0%	658	772.6	91.3%	3	7.4	81.1%

tags in real data recorded by a periodic trigger and, thus, has minimal systematic uncertainty.

Because the atmospheric neutrino background to this search is strongly energy dependent, events are separated into three samples based on visible energy with ranges $100 \text{ MeV} < E_{\text{vis}} < 1.33 \text{ GeV}$, $1.33 \text{ GeV} < E_{\text{vis}} < 20 \text{ GeV}$, and $E_{\text{vis}} > 20 \text{ GeV}$. The number of data events is shown for each sample in Table I, along with the simulated atmospheric neutrino Monte Carlo (MC) expectation. The signal efficiency at representative energies based on signal electron MC simulation is also shown. The importance of the decay electron and neutron tagging cuts is particularly evident in the highest-energy sample ($E_{\text{vis}} > 20 \text{ GeV}$), where they together reduce the background by about a factor of 10 while having a minimal effect on the signal efficiency.

Since boosted dark matter is expected to originate in regions of high dark matter density, this search looks for a signal coming from the Galactic center or the Sun (some boosted dark matter models predict a significant capture rate of cold dark matter in the Sun, either through a spin-dependent dark matter-nucleus cross section [20] or through the combination of a relatively strong dark matter self-interaction and coupling between cold dark matter and standard model particles through boosted dark matter loops [23]). Cones are drawn around the signal source, and the number of events passing the analysis cuts in each cone is counted. While the exact relationship between the boosted dark matter and the scattered electron directions is model dependent, the scattering is, in general, expected to be strongly forward, since the energy of the boosted dark matter and the electron recoil energy are both assumed to be much greater than the mass of the electron. Under these assumptions, the scattering angle of the recoil electron is kinematically constrained to be less than 5.8° at a recoil energy of 100 MeV, with the maximum allowed scattering angle decreasing as $1/\sqrt{E_e}$. Thus, since the reconstructed electron direction was found to have an angular resolution better than 3° over the entire energy range, it is a good proxy for the direction of the boosted dark matter.

When the source of the signal is the Galactic center, the optimal size of the search cone is dependent on both the distribution of the dominant dark matter species in the

Galaxy and the production method of the boosted dark matter. Production of boosted dark matter through both dark matter annihilation and decay were considered for three dark matter halo models: Moore [43], Navarro-Frenk-White (NFW) [44], and Kravtsov [45]. For each combination of halo model and production method, the signal MC was reweighted assuming the direction of the scattered electron was the same as the direction of the boosted dark matter. Optimal cone angles were found by maximizing efficiency/ $\sqrt{\text{background}}$. The optimal half-opening angle of the search cone was found to range from less than 5° to around 40° , depending on the halo model and boosted dark matter production method. We therefore used eight search cones around the Galactic center, ranging from 5° to 40° in steps of 5° . When the Sun is the signal source, the situation is much simpler, since it is effectively a point source. Therefore, a single search cone of 5° around the Sun was used for the solar search.

A data-driven off-source method was used to estimate the background due to atmospheric neutrinos for the two lower-energy samples [46]. In order to avoid contamination from a potential signal, the off-source region is defined as everything outside an 80° cone for the Galactic center search and everything outside of the 5° search cone for the solar search. The off-source regions are defined, like the search cones, in celestial coordinates for the Galactic center search and solar coordinates for the Sun search. For a particular search cone, each data event in the off-source region can be assigned two values based on its direction in horizontal coordinates \hat{d} : T_{off} , which is the fraction of time \hat{d} spends within the off-source region, and T_{on} , which is the fraction of time \hat{d} spends within the search cone. The event is then weighted by the ratio $T_{\text{on}}/T_{\text{off}}$. The sum of these weights gives an estimate of the background in the search cone, while the square root of the sum of the squares of the weights gives the uncertainty on this estimate. The resulting estimates are independent of MC calculations, and their uncertainties range from 3% to 19%, which are smaller than the uncertainties on the corresponding estimates based on MC calculations.

While the off-source method works well for the two lower-energy samples, there are too few events in the highest-energy sample ($E_{\text{vis}} > 20 \text{ GeV}$) for it to be

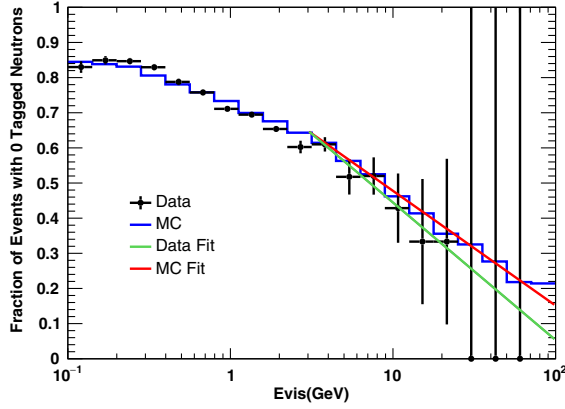


FIG. 2. Neutron tagging cut systematic error estimation. Data and MC fits are to logarithmic function.

successfully applied. Therefore, the background in the highest-energy sample was estimated using the 500-yr SK atmospheric neutrino MC calculation. The MC calculation was livetime normalized and oscillated according to three-flavor oscillations with nominal oscillation parameters. The systematic uncertainty on the estimated background was found by summing in quadrature the effects of 1σ shifts of 75 systematics [40]. The uncertainties in the values of oscillation parameters were included as systematics. The total uncertainty in the background was estimated to be 30%. The dominant systematic is the uncertainty related to the neutron tagging cut, which is estimated to be 23%. This systematic accounts for the uncertainty in the efficiency of the neutron tagging algorithm, as well as in the modeling of production and transport of neutrons in the detector. It was estimated using a data-MC comparison of the fraction of events (over the entire sky) passing the first three analysis cuts that also had zero tagged neutrons. This comparison is shown in Fig. 2. Above about 10 GeV of visible energy, there are very few events in the data, making a data-MC comparison difficult. To compensate, both the data and MC calculations

were fit to logarithmic functions $A + B \log(E_{\text{vis}}/\text{GeV})$, in the region above 3 GeV. The systematic uncertainty as a function of the energy was then taken as the difference between the data fit and the MC fit. The value of the neutron tagging systematic for the background estimate of the highest-energy sample was found by applying this shift on an event by event basis as a function of the visible energy of the particular event. While this estimate is rather imprecise, the expected background rates in the signal cones above 20 GeV are all below one, meaning that the systematic uncertainty on the background estimate is minimal compared to fluctuations associated with Poisson counting statistics. It is thus only the magnitude of this systematic that is important; its exact value has minimal influence on the calculated limits.

The results of the search are shown in Table II. The observed data are consistent with expected background for both the Galactic center and Sun searches. In the highest-energy sample, the search is essentially background-free, and no candidates were found in any of the search cones. Sky maps of the locations of every event passing the analysis cuts are provided for each energy sample in Supplemental Material [47]. For each cone and energy sample, confidence intervals for the observed boosted dark matter event rate were computed using a Poisson χ^2 statistic that incorporates the systematic uncertainty on the background estimate through the pull method [48,49]:

$$\chi^2(s) = \min_{\delta} \left[2 \left(E - \mathcal{O} + \mathcal{O} \ln \frac{\mathcal{O}}{E} \right) + \delta^2 \right], \quad (1)$$

where $E = b(1 + \delta\sigma) + s$, b is the estimated number of background events with systematic uncertainty σ , s is the number of signal events being tested, δ is the systematic pull that is minimized over, and \mathcal{O} is the observed number of events. Note that, for the two lower-energy samples, the background systematic uncertainty σ is due to statistical fluctuations associated with the off-source method, while

TABLE II. Estimated backgrounds, numbers of events in data, and signal event rate limits for each cone and each energy sample. The event rate limits are at the 90% confidence level.

Search cone	$100 \text{ MeV} < E_{\text{vis}} < 1.33 \text{ GeV}$			$1.33 \text{ MeV} < E_{\text{vis}} < 20 \text{ GeV}$			$E_{\text{vis}} > 20 \text{ GeV}$		
	Expected bckg	Data	Sig rate limit (kt yr) ⁻¹	Expected bckg	Data	Sig rate limit (kt yr) ⁻¹	Expected bckg	Data	Sig rate limit (kt yr) ⁻¹
GC 5°	8.4 ± 0.7	5	0.017	1.6 ± 0.3	1	0.018	0.016 ± 0.005	0	0.015
GC 10°	32.0 ± 1.9	24	0.023	6.3 ± 0.84	5	0.026	0.060 ± 0.018	0	0.015
GC 15°	72.5 ± 3.5	69	0.078	13.6 ± 1.6	11	0.032	0.14 ± 0.04	0	0.014
GC 20°	126.5 ± 5.4	125	0.123	23.3 ± 2.3	18	0.028	0.25 ± 0.07	0	0.014
GC 25°	196.8 ± 7.6	202	0.201	35.4 ± 3.3	31	0.049	0.37 ± 0.11	0	0.013
GC 30°	283.7 ± 10.1	285	0.214	49.3 ± 4.3	48	0.081	0.53 ± 0.16	0	0.012
GC 35°	384.8 ± 12.8	375	0.187	68.1 ± 5.4	67	0.101	0.70 ± 0.21	0	0.011
GC 40°	499.6 ± 15.9	494	0.249	90.2 ± 6.9	90	0.124	0.90 ± 0.27	0	0.011
Sun 5°	7.59 ± 0.18	5	0.017	1.25 ± 0.07	1	0.020	0.015 ± 0.004	0	0.015

for the highest-energy sample it is due to uncertainties in the MC production. The test statistic $\Delta\chi^2$ was calculated by subtracting the global minimum χ^2 . To find the confidence level at which a particular value of s is allowed, the measured value of $\Delta\chi^2$ at that value of s was compared to the $\Delta\chi^2$ distribution of a large number of toy MC calculations produced assuming that level of signal. Since the $s = 0$ hypothesis is allowed at 90% confidence for all search cones and energy samples, the upper ends of the 90% confidence intervals are interpreted as 90% upper limits and presented in Table II.

To demonstrate the application of this result to a specific model, limits were calculated on a baseline boosted dark matter model with the Galactic center as the signal source [18]. This model introduces two dark fermions ψ_A and ψ_B and a massive dark photon γ' , with an assumed mass ordering $m_A > m_B > m_{\gamma'}$. The particle ψ_A is proposed to be the dominant cold dark matter in the Universe and does not couple directly to standard model particles. The particle ψ_B is the boosted dark matter and couples to standard model particles through the exchange of the dark photon γ' . The coupling between γ' and ψ_B is set by a coupling constant g' , which is proposed to be large but perturbative, while the coupling between γ' and e^- is scaled from $\gamma - e^-$ coupling by the constant ϵ . Limits were calculated for two scenarios of ψ_B production: one where ψ_B is produced through annihilation of ψ_A with $\bar{\psi}_A$ and another where ψ_B is produced through the decay of ψ_A . In the annihilation scenario, the thermal annihilation cross section is set to $\langle\sigma_{A\bar{A}\rightarrow B\bar{B}}v\rangle = 5 \times 10^{-26} \text{ cm}^3/\text{s}$ in order to achieve the observed relic density $\Omega_A \approx 0.2$ through an assisted freeze-out scenario [18,26]. The energy of ψ_B is equal to m_A in this scenario. In the decay scenario, the decay lifetime of ψ_A , τ_{decay} , is taken to be a free parameter, and the energy of ψ_B is assumed to be $m_A/2$.

Limits were calculated separately for the Moore, NFW, and Kravtsov galactic halo models, using the results from a different cone for each fit. For the annihilation scenario, the 5° cone was used for the Moore model, the 10° cone for the NFW model, and the 40° cone for the Kravtsov model. For the decay scenario, the 40° cone was used for all three galactic halo models. These cones were selected using the cone optimization technique described earlier. For each halo model, signal MC events were reweighted based on the values of m_A , m_B , $m_{\gamma'}$, ϵ , and g' at the particular point in parameter space being tested. This reweighting accounts for the model-dependent recoil electron energy spectrum, as well as the model-dependent smearing between the boosted dark matter direction and the recoil electron direction. The effect of boosted dark matter scattering off of both electrons and protons in Earth is also accounted for, though this effect is negligible for the majority of the allowed parameter space. A binned χ^2 statistic was then computed similar to the one described above:

$$\chi^2 = \sum_i^3 \min_{\delta_i} \left[2 \left(E_i - O_i + O_i \ln \frac{O_i}{E_i} \right) + \delta_i^2 \right], \quad (2)$$

with variables defined as before, summed over three bins corresponding to the three energy samples. The $\Delta\chi^2$ test statistic was then calculated by subtracting the global minimum χ^2 . Confidence intervals were found by comparing the measured $\Delta\chi^2$ values with the distributions of $\Delta\chi^2$ values found by many toy Monte Carlo simulations produced at each point. Ninety-percent confidence intervals were computed in the ϵ vs m_A plane for the annihilation scenario, and the $\epsilon/\tau_{\text{decay}}$ vs m_A plane for the decay scenario, with m_B , $m_{\gamma'}$, and g' set to representative values of $m_B = 200 \text{ MeV}$, $m_{\gamma'} = 20 \text{ MeV}$, and $g' = 0.5$, respectively. Since the $\epsilon = 0$ points, which correspond to no signal, are allowed at 90% confidence, the resulting confidence intervals are interpreted as upper limits. These limits are shown for the Moore, NFW, and Kravtsov halo models in Fig. 3.

In summary, we have searched for evidence of boosted dark matter by looking for high-energy elastically scattered electrons that point back to the Galactic center or the Sun. We have found no such evidence. This is the first study of high-energy electron elastic scatterlike events at SK. The use of decay electron and tagged neutron cuts significantly reduced background in the highest-energy sample, allowing

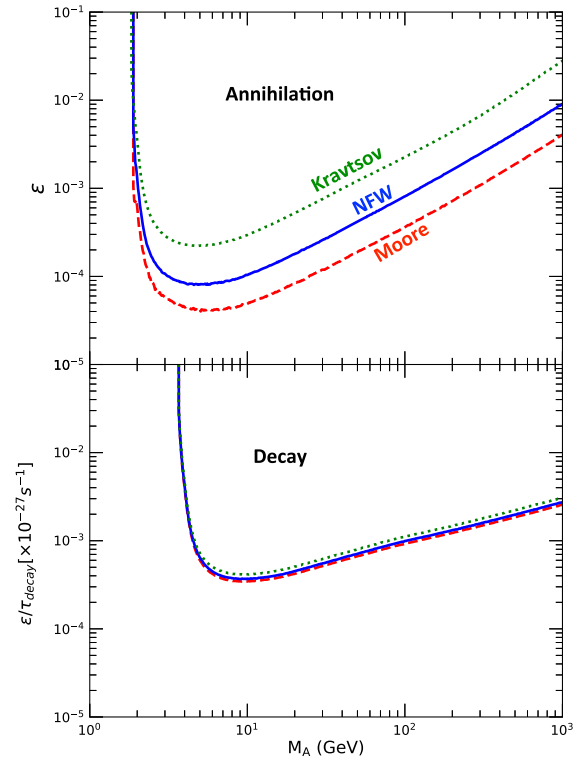


FIG. 3. 90% confidence interval upper limits for $m_B = 200 \text{ MeV}$, $m_{\gamma'} = 20 \text{ MeV}$, and $g' = 0.5$, for boosted dark matter produced by annihilation (top) and decay (bottom).

for an effectively background-free search in that energy range. Our results are presented in a model-independent way, which makes them applicable not only to boosted dark matter but to any theory that predicts an excess of particles from the Galactic center or Sun that would elastically scatter electrons to energies above 100 MeV.

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